

ELECTRICAL AND MAGNETIC PROPERTIES OF MULTIFERROIC $\text{Co}_{0.6}\text{Zn}_{0.4}\text{Fe}_{1.7}\text{Mn}_{0.3}\text{O}_4$ DOPED $0.99\text{Bi}_{0.47}\text{Na}_{0.47}\text{Ba}_{0.06}\text{TiO}_3$ - $0.01\text{Ba}(\text{Sn}_{0.70}\text{Nb}_{0.24})\text{O}_3$ CERAMICS SYNTHESIZED VIA THE SOLID-STATE COMBUSTION TECHNIQUE

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Abstract

Multiferroic $(1-x)[0.99\text{Bi}_{0.47}\text{Na}_{0.47}\text{Ba}_{0.06}\text{TiO}_3-0.01\text{Ba}(\text{Sn}_{0.70}\text{Nb}_{0.24})\text{O}_3]-x\text{Co}_{0.6}\text{Zn}_{0.4}\text{Fe}_{1.7}\text{Mn}_{0.3}\text{O}_4$ (abbreviated as BNBT-BSN- x CZFMO) ceramics with x ranging from 0 to 0.20 were fabricated using the solid-state combustion technique. The effect of varying BNBT-BSN : CZFMO ratios on the phase structure, microstructure, electrical and magnetic properties was investigated. X-ray diffraction (XRD) analysis of pure BNBT-BSN showed a perovskite structure with rhombohedral and tetragonal phases. The doped BNBT-BSN- x CZFMO ceramics displayed coexisting rhombohedral, tetragonal, and cubic spinel phases, with the cubic spinel phase increasing when the CZFMO content increased. In addition, the XRD peaks shifted to higher angles as the CZFMO content increased, indicating a decrease in lattice parameters. The dielectric constant decreased with higher CZFMO content and higher frequencies. The pure BNBT-BSN ceramic exhibited a saturated P-E loop with a P_{max} of $33.2 \mu\text{C}/\text{cm}^2$, P_r of $26.1 \mu\text{C}/\text{cm}^2$, and an E_c of $14.5 \text{ kV}/\text{cm}$. With increased CZFMO content, non-saturated and bloated P-E loops with lower P_{max} , P_r , and E_c were observed, implying a rise in the leakage current. The addition of CZFMO induced ferromagnetic behavior in the ceramics, leading to an increase in M_s and a reduction in H_c as CZFMO content increased. The magnetoelectric coupling coefficient of BNBT-BSN- x CZFMO ceramics continuously increased with higher CZFMO content.

Keywords: BNBT-BSN- x CZFMO; Dielectric; Ferroelectric; Ferromagnetic; Phase Structure

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Introduction

Due to their excellent electrical and piezoelectric properties, the electronic device industry commonly uses lead-based ceramics, such as lead zirconate titanate (PZT). However, the lead vapors produced during synthesis are hazardous to human health and the environment. Consequently, there is a need to replace lead-based ceramics with lead-free alternatives like bismuth sodium titanate ($\text{Bi}_{0.5}\text{Na}_{0.5}\text{TiO}_3$), potassium sodium niobate ($\text{K}_{0.5}\text{Na}_{0.5}\text{NbO}_3$), and barium titanate (BaTiO_3) (Li *et al.*, 2013; Parija *et al.*, 2013; Lee *et al.*, 2015). Bismuth sodium titanate ($\text{Bi}_{0.5}\text{Na}_{0.5}\text{TiO}_3$; BNT) lead-free ceramics are particularly attractive due to their superior electrical properties, including a high Curie temperature ($T_C=540^\circ\text{C}$) and remnant polarization ($P_r=38 \mu\text{C}/\text{cm}^2$). However, a large coercive field ($E_C=73 \text{ kV}/\text{cm}$) limits practical application, resulting in a low piezoelectric coefficient (d_{33}) (Nagata *et al.*, 2006). To fix these limitations, solid solutions such as barium titanate (BaTiO_3), barium calcium titanate ($\text{Ba}_{0.7}\text{Ca}_{0.3}\text{TiO}_3$), bismuth potassium titanate ($\text{Bi}_{0.5}\text{K}_{0.5}\text{TiO}_3$), and strontium titanate (SrTiO_3) have been added to BNT ceramics (Zou *et al.*, 2010; Chandrasekhar and Kumar, 2015; Cao *et al.*, 2016; Wang and Shi, 2023). Introducing 6% BaTiO_3 into the BNT system produced $\text{Bi}_{0.47}\text{Na}_{0.47}\text{Ba}_{0.06}\text{TiO}_3$ (BNBT) ceramics at a morphotropic phase boundary (MPB) condition. This composition, recently fabricated by B. Thatawong *et al.* (Thatawong *et al.*, 2022) using the solid-state combustion technique, showed a high maximum dielectric constant ($\epsilon_m=8405$), remnant polarization ($P_r=28.2 \mu\text{C}/\text{cm}^2$), and piezoelectric coefficient ($d_{33}=161 \text{ pC}/\text{N}$). The structure's complexity allows for modification using various conditions like electric field and temperature, which can enhance the ceramic's electrical properties (Li *et al.*, 2016).

The BNBT system at the MPB requires modifications to optimize its properties further. This generally involves adding a solid solution with a similar structure to create a ternary system in BNBT, enhancing specific properties like dielectric and ferroelectric performance (Swain *et al.*, 2015). Y. Jia *et al.* (Jia *et al.*, 2015) incorporated $\text{Ba}(\text{Sn}_{0.70}\text{Nb}_{0.24})\text{O}_3$ (BSN) into BNBT using the traditional solid-state method, resulting in high strain properties (S_{max} of 0.43%, d_{33}^* of 633 pm/V), a good maximum dielectric constant (ϵ_m) of approximately 4000 and a large maximum polarization (P_{max}) of around $40 \mu\text{C}/\text{cm}^2$. Recently, A. Luangpangai *et al.* (Luangpangai *et al.*, 2024) successfully fabricated $0.99\text{Bi}_{0.47}\text{Na}_{0.47}\text{Ba}_{0.06}\text{TiO}_3$ - $0.01\text{BaSn}_{0.70}\text{Nb}_{0.24}\text{O}_3$ (BNBT-BSN) ceramics using

the solid-state combustion technique. The ceramics were calcined at 800°C for 2 hours and sintered at 1175°C for 2 hours. This ceramic exhibited an MPB between rhombohedral and tetragonal phases with an excellent ϵ_m of 6513, a large P_r of $28.98 \mu\text{C}/\text{cm}^2$, and a high P_{max} of $42.21 \mu\text{C}/\text{cm}^2$.

Ferromagnetic materials with a spinel ferrite structure (AB_2O_4) have been widely employed in various devices, with potential future applications in multiferroic systems (Xiao *et al.*, 2010). Among these, a $\text{Co}_{0.6}\text{Zn}_{0.4}\text{Fe}_{1.7}\text{Mn}_{0.3}\text{O}_4$ (CZFMO) ceramic was successfully synthesized by P. Thawong *et al.* (Thawong *et al.*, 2019) using the solid-state combustion technique with a low firing temperature and short dwell time (calcined at 750°C for 2 hours and sintered at 1200°C for 2 hours). The ceramic showed a good dielectric constant (ϵ) of 365 and high ferromagnetic properties (M_s of $72.50 \text{ emu}/\text{g}$ and H_c of 4.24 Oe).

Multiferroic materials are typically made by combining two or more ferroelectric, ferromagnetic, or ferroelastic phases. The goal of the combination is to take advantage of the specific properties of each phase and improve the overall multiferroic behavior. Many reports have documented progress in developing lead-free multiferroic materials that feature a perovskite structure with ferroelectric and spinel ferrite ferromagnetic phases. For example, CZFMO-modified $\text{Ba}_{0.85}\text{Ca}_{0.15}\text{Ti}_{0.9}\text{Zr}_{0.1}\text{O}_3$ systems exhibited improved ferromagnetic properties (M_s of $6.32 \text{ emu}/\text{g}$ and M_r of $0.55 \text{ emu}/\text{g}$) and a magnetoelectric coupling coefficient (α_E) of $-2.415 \text{ mV}/\text{Oe}\cdot\text{cm}$, as reported by M. Singh *et al.* (Singh *et al.*, 2020). Similarly, E. Venkata Ramana *et al.* (Venkata Ramana *et al.*, 2012) prepared BNBT ceramics doped with CZFMO using the conventional solid-state reaction technique. This ceramic demonstrated a good P_r of $12.5 \mu\text{C}/\text{cm}^2$, an M_s of $13.6 \text{ emu}/\text{g}$, and an α_E of $7.8 \text{ mV}/\text{Oe}\cdot\text{cm}$. This suggests that adding CZFMO, with its spinel structure, into lead-free ceramics is a promising approach for improving ferroelectric, ferromagnetic, and magnetoelectric properties. As a result, the potential of improving these properties in BNBT-BSN with CZFMO is interesting. Notably, the investigation of multiferroic BNBT-BSN-CZFMO ceramics remains largely unexplored.

In this study, multiferroic $(1-x)[0.99\text{Bi}_{0.47}\text{Na}_{0.47}\text{Ba}_{0.06}\text{TiO}_3-0.01\text{Ba}(\text{Sn}_{0.70}\text{Nb}_{0.24})\text{O}_3]-x\text{Co}_{0.6}\text{Zn}_{0.4}\text{Fe}_{1.7}\text{Mn}_{0.3}\text{O}_4$ (BNBT-BSN- x CZFMO) ceramics with x between 0 and 0.20 were synthesized using the solid-state combustion technique and their phase structure, microstructure, electrical, magnetic and magnetoelectric properties were investigated.

Materials and Methods

The solid-state combustion technique was used to fabricate multiferroic

$(1-x)[0.99\text{Bi}_{0.47}\text{Na}_{0.47}\text{Ba}_{0.06}\text{TiO}_3 - 0.01\text{Ba}(\text{Sn}_{0.70}\text{Nb}_{0.24})\text{O}_3] - x\text{Co}_{0.6}\text{Zn}_{0.4}\text{Fe}_{1.7}\text{Mn}_{0.3}\text{O}_4$ (BNBT-BSN- x CZFMO) ceramics with x ranging from 0 to 0.20. The BNBT-BSN and CZFMO components were prepared separately using the stoichiometric ratio. $\text{Bi}(\text{NO}_3)_3 \cdot 5\text{H}_2\text{O}$ (98.50%, KEMAUS, AU), NaNO_3 (99.00%, Ajex Finechem, AU), $\text{Ba}(\text{NO}_3)_2$ (99.00%, HiMedia, USA), TiO_2 (99.00%, KEMAUS, AU), SnO_2 (99.00%, Sigma-Aldrich, USA) and Nb_2O_5 (99.00%, Eleps, RUS) were the raw materials for BNBT-BSN. While, CZFMO raw materials consisted of $\text{Co}(\text{NO}_3)_2 \cdot 6(\text{H}_2\text{O})$ (98.00%, Ajex Finechem, AU), ZnO (99.00%, Merck German, DE), $\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$ (101.00%, Ajex Finechem, AU) and MnO_2 (99.00%, Sigma-Aldrich, USA). The raw materials were ball milled using zirconium balls in an ethanol solution for 24 hours. After drying, the slurries were ground and combined with glycine ($\text{C}_2\text{H}_5\text{NO}_2$: 99.00%, KemAus, AUS) as fuel at a propellant ratio of 1.1111 wt% for BNBT-BSN and 0.855 wt% for CZFMO. The BNBT-BSN and CZFMO powders were calcined separately at 850°C for 2 hours and 750°C for 2 hours, respectively, with a 5°C/min heating rate. The calcined BNBT-BSN and CZFMO powders were combined at the desired weight ratio ($x = 0$ to 0.20), mixed with polyvinyl butyral (PVB), and pressed into green pellets with a diameter of 15 mm and a thickness of 1 mm under 80 MPa of pressure. Finally, the green pellets were sintered at 1175°C for 2 hours, with a heating rate of 2°C/min.

The phase structure of the ceramics was characterized by X-ray diffraction (XRD, Bruker, D8 Discover) with $\text{Cu-K}\alpha 1,2$ radiation, scanning in the 2θ range of 10° to 70° at room temperature. Phase analysis was conducted using the Rietveld refinement method via the Fullprof program. The microstructure of the sample surfaces was examined with scanning electron microscopy (SEM, Thermo Fisher Scientific, Apreo S), after thermally etching 100°C below the sintering temperature. The average grain size was determined using the linear intercept method, based on measurements of over 300 grains. Sample density was calculated using Archimedes method.

To measure the electrical and magnetoelectric properties, silver paste (IRMS sp5574, CN) was coated on both sides of the samples and fired at 550°C for 20 min to make electrodes. Temperature and frequency-dependent dielectric properties were measured using LCR meters (Agilent 4284A and

HP4284A, respectively). The polarization versus electric field (P-E) hysteresis loops were measured with a computer-controlled modified Sawyer-Tower circuit (Radiant, PLC2-The P-E loop). Ferromagnetic hysteresis (M-H) loops were obtained using a vibrating sample magnetometer (Versa Lab, Quantum Design). Magnetoelectric coupling was analyzed with a lock-in amplifier (Stanford Research System, model SR830).

Results and Discussion

X-ray diffraction (XRD) was used to determine the phase structure of BNBT-BSN- x CZFMO ceramics. Figure 1(a) displays the XRD patterns of the BNBT-BSN- x CZFMO samples with CZFMO content between 0 and 0.20, in the 2θ range of 10 to 70°. The pure BNBT-BSN ceramic showed a pure perovskite structure without any secondary phases. As the CZFMO content increased, the coexistence of perovskite and spinel phases was observed. Moreover, secondary phases of Bi_2O_3 (JCPDS files No. 29-0236) and Fe_2O_3 (JCPDS files No. 33-0664) were found for samples with $x = 0.10, 0.15$ and 0.20.

To further analyze the phase structure, Figure 1(b) focuses on the XRD in the 2θ range of 38° to 50°. This region includes indexed peaks for both perovskite and spinel phases. According to JCPDS data No. 36-0339 (Sinkruason *et al.*, 2023), the tetragonal phase is identified by a single $(111)_T$ peak at 2θ of 40° and dual $(002)_T/(200)_T$ peaks at 2θ of 46°. While the rhombohedral phase is characterized by dual $(003)_R/(021)_R$ peaks at 2θ of 40° and a single $(202)_R$ peak at 2θ of 46°, according to JCPDS files No. 36-0340 (Sinkruason *et al.* 2023). Also, the cubic spinel phase is indexed by single $(400)_C$ and $(331)_C$ peaks at $2\theta = 43^\circ$ and 47° , respectively, according to JCPDS files No. 22-1086 (Jiang and Ai, 2010).

Pure BNBT-BSN ceramics showed split peaks at 2θ of 40° and 46°, indicating the coexistence of tetragonal and rhombohedral phases, consistent with observations by A. Luangpangai *et al.* (Luangpangai *et al.*, 2024). When CZFMO content increased, these split peaks merged into asymmetric peaks. A small distinct single peak emerged at higher amounts of CZFMO ($x = 0.15$ and 0.20), suggesting the coexistence of perovskite and spinel phases. In addition, XRD peaks at 2θ of 40° and 46° moved slightly to a lower angle, implying an expansion in lattice parameters. This shift may be attributed to the substitution of Ti (0.605 Å) and Nb (0.640 Å) ions with larger Fe (0.645 Å) and Mn (0.645 Å) ions at the B-sites, leading to an extended lattice structure.

To quantify the phase structure content in BNBT-BSN- x CZFMO ceramics, the Rietveld refinement method was applied using the Fullprof program, as seen in Figure 2(a-e). Data sources for determining the lattice parameters, including length (a , b , c), angle (α , β , γ), atomic data, and occupancy

rate, were obtained from the Crystallography Open Database (COD). To fit the XRD results of BNBT-BSN- x CZFMO ceramics, a combination of structures including tetragonal (P4bm space group), rhombohedral (R3c space group) and cubic spinel phase (Fd3m space group) was applied. A pseudo-

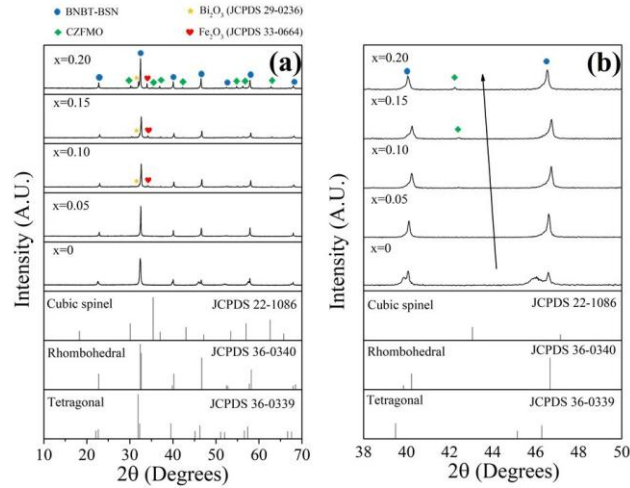


Figure 1. XRD patterns of BNBT-BSN- x CZFMO ceramics with different x , showing phase evolution with increasing CZFMO content, highlighting perovskite and spinel phases, at the 2θ range of (a) 10-70° and (b) 38-50°

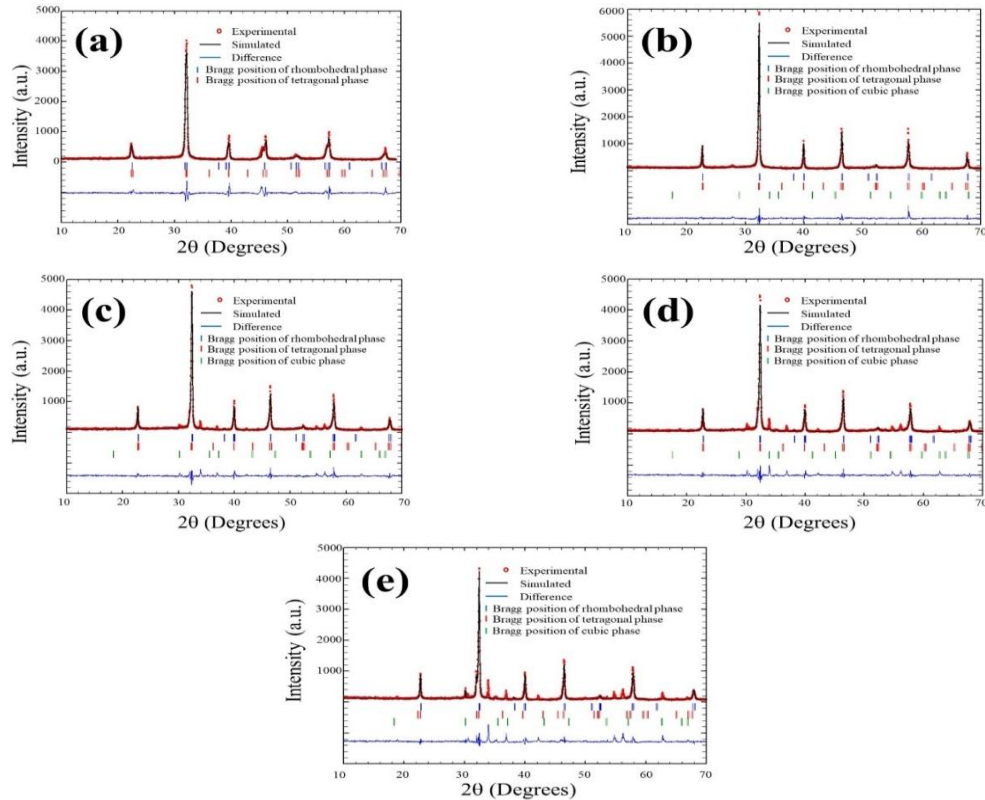


Figure 2. Rietveld refinement results for BNBT-BSN- x CZFMO ceramics with (a) $x=0$, (b) $x=0.05$, (c) $x=0.10$, (d) $x=0.15$, and (e) $x=0.20$

Voigt function was used to model the XRD peaks, while a Chebyshev polynomial fit the background. The fitting results, including lattice parameters and phase percentages, are presented in Table 1. The undoped BNBT-BSN ceramic showed a coexistence of tetragonal (T) and rhombohedral (R) phases with a nearly equal T-to-R phase ratio of 55:45, consistent with previous reports (Luangpangai *et al.*,

2024). However, a cubic spinel (C) phase emerged alongside the R and T phases when CZFMO content increased. The proportion of the C phase continuously increased from 2% ($x=0.05$) to 7% ($x=0.10$), 14% ($x=0.15$), and 35% ($x=0.20$), as shown in Table 1. The increase in the C phase corresponds with the rising CZFMO content.

Table 1. The fitting parameters, lattices parameters, and phase percentages of BNBT-BSN- x CZFMO ceramics with $x=0-0.20$

Samples	Refine parameters	Phase structures	Profile parameters	Atoms information					The percentage of phase
				Label	x	y	z	Occ.	
$x=0$	$\chi^2=3.02$ $R_p=28.3\%$ $R_{wp}=26.1\%$ $R_{exp}=15.0\%$	$R3c:H$	$a=5.5667\text{\AA}$ $c=13.8347\text{\AA}$ $u=0.2620$ $v=0.3810$ $w=0.0085$	Na	0	0	0.2627	0.37280	45
				Bi	0	0	0.2627	0.38668	
				Ba	0	0	0.2627	0.06650	
				Ti	0	0	0.0063	0.66690	
				Sn	0	0	0.0063	0.04982	
				Nb	0	0	0.0063	0.01996	
				O	0.1260	0.3360	0.0833	3.23624	
		$P4bm$	$a=5.5607\text{\AA}$ $c=3.9799\text{\AA}$ $u=0.4873$ $v=0.3131$ $w=0.0889$	Na	0	0.5	0.5450	0.41677	55
				Bi	0	0.5	0.5450	0.53575	
				Ba	0	0.5	0.5450	0.07168	
				Ti	0	0	0	0.93094	
				Sn	0	0	0	0.01012	
				Nb	0	0	0	0.15076	
$x=0.05$	$\chi^2=1.56$ $R_p=23.7\%$ $R_{wp}=20.6\%$ $R_{exp}=16.52\%$	$R3c:H$	$a=5.5278\text{\AA}$ $c=13.5604\text{\AA}$ $u=0.1211$ $v=0.1501$ $w=0.0211$	O1	0	0	0.2627	0.22182	35
				Bi	0	0	0.2627	0.38644	
				Ba	0	0	0.2627	0.03474	
				Ti	0	0	0.0063	0.61860	
				Sn	0	0	0.0063	0.03766	
				Nb	0	0	0.0063	0.01921	
				O	0.1260	0.3360	0.0833	3.39273	
		$P4bm$	$a=5.5512\text{\AA}$ $c=3.9011\text{\AA}$ $u=1.4977$ $v=0.1387$ $w=0.0012$	Na	0	0.5	0.5450	0.55868	63
				Bi	0	0.5	0.5450	0.53618	
				Ba	0	0.5	0.5450	0.07593	
				Ti	0	0	0	0.94505	
				Sn	0	0	0	0.01222	
				Nb	0	0	0	0.15113	
				O1	0	0	0.5100	0.90793	2
				O2	0.2710	0.2290	0.0150	1.76005	
		$Fd-3m$	$a=b=c=8.7308\text{\AA}$ $u=2.7826$ $v=2.6580$ $w=0.2997$	O1	0.25000	0.25000	0.25000	1.41573	
				Fe1	0.12500	0.12500	0.12500	0.28913	
				Fe2	0.50000	0.50000	0.50000	0.25732	
				Mn1	0.12500	0.12500	0.12500	0.04996	
				Mn2	0.50000	0.50000	0.50000	0.09905	
				Zn1	0.12500	0.12500	0.12500	0.18479	
$x=0.10$	$\chi^2=2.21$ $R_p=25.8\%$ $R_{wp}=24.4\%$ $R_{exp}=16.45\%$	$R3c:H$	$a=5.5183\text{\AA}$ $c=13.5696\text{\AA}$ $u=0.1390$ $v=0.1842$ $w=0.0193$	Zn2	0.50000	0.50000	0.50000	0.37077	33
				Co1	0.12500	0.12500	0.12500	0.19708	
				Co2	0.50000	0.50000	0.50000	0.15199	
				Na	0	0	0.2627	0.69141	
				Bi	0	0	0.2627	0.53550	
				Ba	0	0	0.2627	0.07603	
		$P4bm$	$a=5.5474\text{\AA}$ $c=3.8945\text{\AA}$ $u=0.2809$ $v=0.3281$ $w=0.1429$	Ti	0	0	0.0063	0.93320	60
				Sn	0	0	0	0.01244	
				Nb	0	0	0	0.15101	
				O1	0	0	0.5100	1.10965	
				O2	0.2710	0.2290	0.0150	2.13461	
				O1	0.25000	0.2290	0.25000	0.96274	7
		$Fd-3m$	$a=b=c=8.3805\text{\AA}$ $u=4.8083$ $v=1.0261$ $w=2.6383$	Fe1	0.12500	0.25000	0.12500	0.25924	
				Fe2	0.50000	0.12500	0.50000	0.44891	
				Mn1	0.12500	0.50000	0.12500	0.06639	
				Mn2	0.50000	0.12500	0.50000	0.16551	
				Zn1	0.12500	0.50000	0.12500	0.20346	
				Zn2	0.50000	0.12500	0.50000	0.41741	
				Co1	0.12500	0.50000	0.12500	0.21227	7
				Co2	0.50000	0.12500	0.50000	0.20206	

Table 1. The fitting parameters, lattices parameters, and phase percentages of BNBT-BSN-xCZFMO ceramics with x=0-0.20

Samples	Refine parameters	Phase structures	Profile parameters	Atoms information					The percentage of phase
				Label	x	y	z	Occ.	
x=0.05	$\chi^2=3.03$ $R_p=28.9\%$ $R_{wp}=29.1\%$ $R_{exp}=16.72\%$	R3c:H	a=5.5189Å	Na	0	0	0	0.25054	32
			c=13.5778Å	Bi	0	0	0	0.38575	
			u=−0.1396	Ba	0	0	0	0.03497	
			v=0.1955	Ti	0	0	0	0.53010	
			w=−0.0211	Sn	0	0	0	0.03956	
				Nb	0	0	0	0.01905	
				O	0.1260	0.3360	0.3360	3.23577	
				Na	0	0	0.5450	0.5450	
				Bi	0	0	0.5450	0.5450	
		P4bm	a=5.5463Å	Ba	0	0	0.5450	0.5450	54
			c=3.9015Å	Ti	0	0	0	0	
			u=−0.2404	Sn	0	0	0	0	
			v=0.6001	Nb	0	0	0	0	
			w=0.1616	O1	0	0	0.5100	0.5100	
				O2	0.2710	0.2710	0.0150	0.0150	
				O1	0.25000	0.25000	0.25000	0.34779	
				Fe1	0.12500	0.12500	0.12500	0.33121	
				Fe2	0.50000	0.50000	0.50000	0.16784	
		Fd-3 m	a=b=c=8.7684Å	Mn1	0.12500	0.12500	0.12500	0.27678	14
			u=7.2525	Mn2	0.50000	0.50000	0.50000	0.07761	
			v=8.1273	Zn1	0.12500	0.12500	0.12500	0.20674	
			w=−0.4732	Zn2	0.50000	0.50000	0.50000	0.36016	
				Co1	0.12500	0.12500	0.12500	0.17308	
				Co2	0.50000	0.50000	0.50000	0.15212	
				Na	0	0	0.2627	0.34239	
				Bi	0	0	0.2627	0.38720	
				Ba	0	0	0.2627	0.03474	
R3c:H	a=5.5139Å	Ti	0	0	0.0063	0.60880	45		
	c=13.5570Å	Sn	0	0	0.0063	0.03943			
	u=−0.1077	Nb	0	0	0.0063	0.01945			
	v=0.1759	O	0.1260	0.3360	0.0833	5.56031			
	w=−0.0271	Na	0	0.5	0.5450	0.16208			
		Bi	0	0.5	0.5450	0.70116			
		Ba	0	0.5	0.5450	0.07574			
		Ti	0	0	0	0.96739			
		Sn	0	0	0	0.01247			
P4bm	a=5.5407Å	Nb	0	0	0	0.15058	20		
	c=3.9926Å	O1	0	0	0.5100	0.37308			
	u=1.0289	O2	0.2710	0.2290	0.0150	6.21947			
	v=−0.2569	O1	0.25000	0.25000	0.25000	1.38011			
	w=0.0314	Fe1	0.12500	0.12500	0.12500	0.08221			
		Fe2	0.50000	0.50000	0.50000	0.30418			
		Mn1	0.12500	0.12500	0.12500	0.12982			
		Mn2	0.50000	0.50000	0.50000	0.17748			
		Zn1	0.12500	0.12500	0.12500	0.32256			
Fd-3 m	a=b=c=8.3920Å	Zn2	0.50000	0.50000	0.50000	0.42124	35		
	u=83.4376	Co1	0.12500	0.12500	0.12500	0.53911			
	v=4.3107	Co2	0.50000	0.50000	0.50000	0.53198			
	w=−7.2226								

A scanning electron microscope (SEM) was used to examine the microstructure of BNBT-BSN-xCZFMO ceramics with various CZFMO content. The surface morphology of these ceramics is illustrated in Figure 3(a)-(e). The undoped BNBT-BSN ceramic (x=0) had polygonal grain shapes with anisotropic grain growth (Figure 3(a)). When CZFMO content increased up to x=0.05, three different grain types emerged: rough-surface, smooth-surface, and rod-like grains, as seen in Figure 3(b). As x increased to 0.10 and 0.15, the rough-surface and rod-like grains diminished, while smooth-surface grains increased and merged, as shown in Figures 3(c) and (d). At higher CZFMO content (x=0.20), the smooth-surfaced grains were primarily distributed at the grain boundaries of rough-surface grains, and rod-like grains were not observed, as seen in Figure 3(e). Additionally,

increased CZFMO content increased porosity at the grain boundaries, likely caused by excessive oxygen vacancies (Bongkarn *et al.*, 2016).

Table 2 shows that the pure BNBT-BSN ceramic had an average grain size of 1.58 μm . When CZFMO content increased to 0.05, the average grain size significantly increased to 13.59 μm . This increase could be attributed to forming a liquid phase during the sintering process and creating oxygen vacancies, which enhanced mass transfer when CZFMO content increased (Zhang *et al.*, 2017). However, the average grain size continuously decreased with further increases in CZFMO contents (x>0.05). This suggests that higher amounts of CZFMO accumulated at the grain boundaries, inhibiting grain growth (Kornphom *et al.*, 2023).

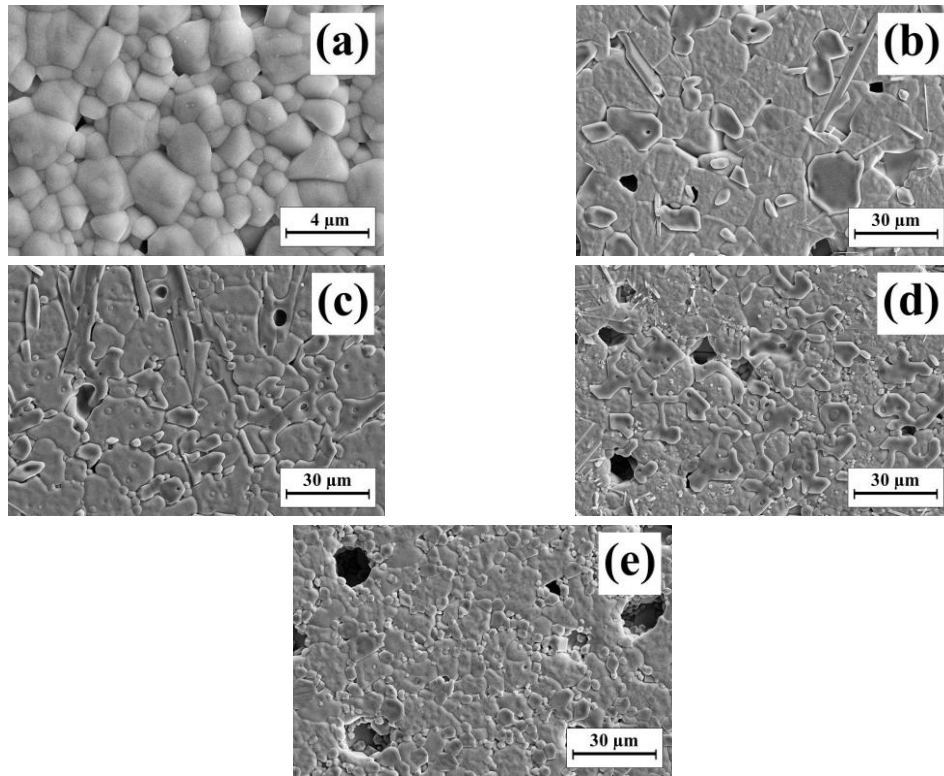


Figure 3. Surface morphology of BNBT-BSN-xCZFMo ceramics with (a) $x=0$, (b) $x=0.05$, (c) $x=0.10$, (d) $x=0.15$, and (e) $x=0.20$

Table 2. The average grain size, density, dielectric, ferroelectric, and ferromagnetic properties of BNBT-BSN-xCZFMo ceramics with various x levels

CZFMo contents	Average grain size (μm)	Density (g/cm^3)	ϵ_r at 1 kHz	$\tan\delta_r$ at 1 kHz	P_{max} ($\mu\text{C}/\text{cm}^2$)	P_r ($\mu\text{C}/\text{cm}^2$)	E_c (kV/cm)	M_s (emu/g)	M_r (emu/g)	H_c (Oe)	α_E (mV/Oe.cm)
$x=0$	1.58	5.54	1505	0.05	33.19	26.09	14.22	-0.0044	9×10^{-5}	52	1.317
$x=0.05$	13.59	5.32	952	0.02	14.68	5.16	11.06	0.0715	9×10^{-4}	24	2.743
$x=0.10$	12.39	5.21	674	0.03	11.30	11.73	24.99	0.1411	9×10^{-5}	7	3.655
$x=0.15$	6.70	5.10	501	0.03	5.45	3.01	17.40	0.2485	7×10^{-5}	3	4.402
$x=0.20$	4.52	5.05	364	0.03	3.80	2.09	17.39	0.8056	-	5	5.111

To investigate each specific grain type in the SEM images, energy dispersive spectroscopy (EDS) was used to identify and quantify the elemental composition of the grains, as shown in Figure 4. In the BNBT-BSN-0.05CZFMo ceramic, the rough-surface, smooth-surface, and rod-like grains were marked and labeled with Spectrum 1, Spectrum 2, and Spectrum 3, respectively. The rough-surface grain had a high concentration of Bi (48.5 wt%) and Ti (21.7 wt%). However, the smooth-surface grain had elevated levels of Co (3.1 wt%), Mn (1.4 wt%) and Fe (7.4 wt%), indicating the presence of ferroelectric and ferrite phases in the respective grains. Additionally, the rod-like grains contained a balance of Bi (13.2 wt%), Ba (13.7 wt%), Ti (32.9 wt%), Co (2.5 wt%), Mn (1.1 wt%) and Fe (4.8 wt%), indicating a coexistence of ferroelectric and ferrite structures, with more ferrite structures.

The EDS mapping confirmed the distribution of elements in the surface morphology of the BNBT-BSN-0.05CZFMo ceramic, as shown in Figure 5. It showed that Bi and Na were predominantly in ferroelectric grains, while only Ba, Ti, Co, and Fe are clearly seen in the ferrite and rod-like grains. This distribution supports the classification of different grain types within the material.

The density of the BNBT-BSN-xCZFMo ceramics gradually decreased from 5.54 to 5.05 g/cm^3 as the CZFMo content increased from 0 to 0.20, as listed in Table 2. This result corresponds with the SEM images, which revealed increased porosity and poor microstructure at higher CZFMo content.

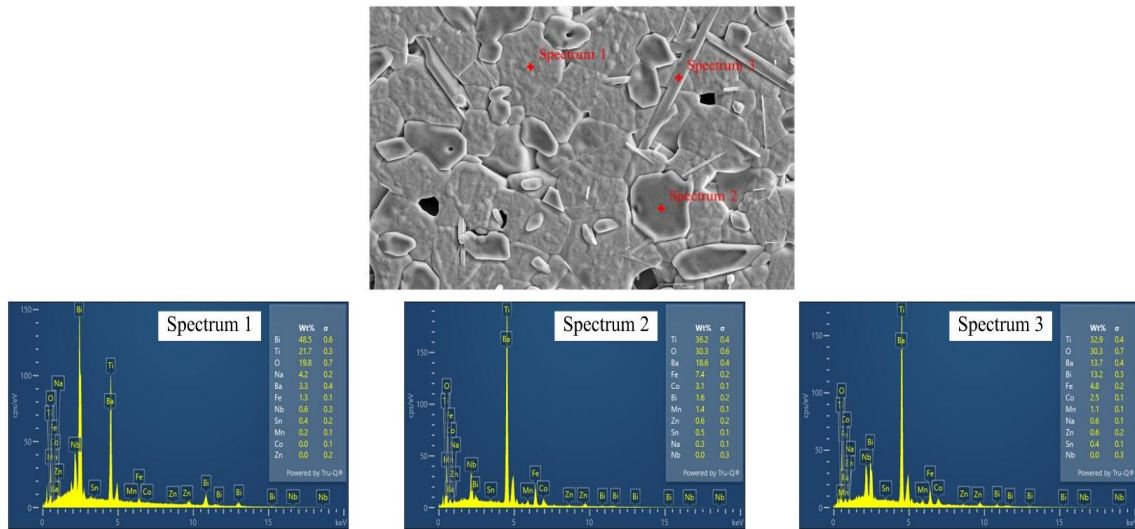


Figure 4. EDS element analysis of the rough surface (spectrum 1), smooth surface (spectrum 2), and rod-like grains (spectrum 3) in the 0.95BNBT-BSN-0.05CZFMO ceramic

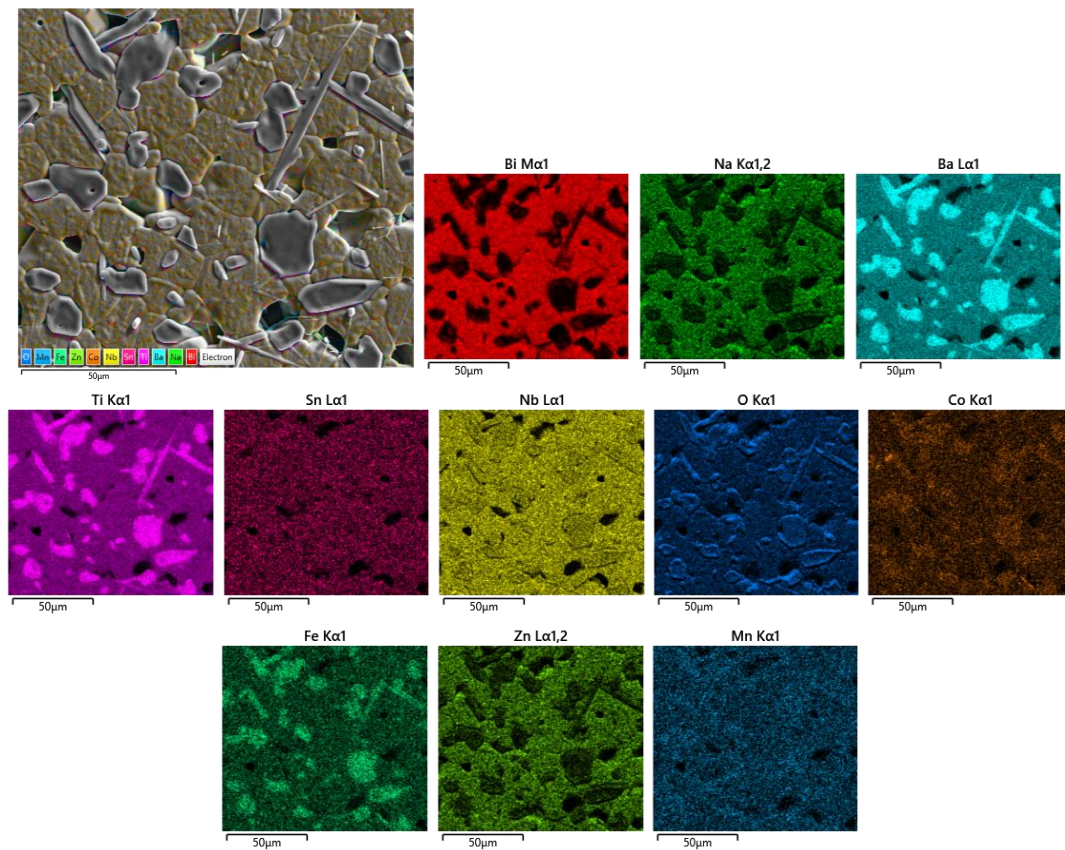


Figure 5. EDS mapping analysis showing the elemental distribution in the surface morphology of the 0.95BNBT-BSN-0.05CZFMO ceramic

The temperature-dependent dielectric constant (ϵ) and loss ($\tan\delta$) of BNBT-BSN-xCZFM0 ceramics with varying CZFMO content are displayed in Figure 6(a-j). The measurements were taken from room temperature to 400°C at frequencies of 1, 10, and 100 kHz. In general, the dielectric curve of BNT-based ceramics shows two dielectric peaks. The first dielectric peak appears at the depolarization temperature (T_d), which marks the thermal transition of polar nanoregions (PNRs) from a rhombohedral structure (R3c symmetry) to a tetragonal structure (P4bm symmetry) (Badapanda *et al.*, 2017). While the second dielectric peak is at T_m , the temperature where the dielectric constant reaches its maximum (Badapanda *et al.*, 2017).

In Figure 6(a), the dielectric curve of the pure BNBT-BSN sample exhibited two dielectric peaks with T_d at 147°C and T_m at 278°C, measured at 1 kHz. The dielectric response is broad with frequency dispersion, characteristic of relaxor ferroelectric behavior, where T_d shifts to higher temperatures and T_m shifts to lower temperatures with increased frequency, as shown in Figure 6(a) (Luangpangai *et al.*, 2024).

When a small amount of CZFMO ($x=0.05$) is added, the dielectric response and loss measured at 1 kHz significantly increased at high temperatures, and T_m disappears, as seen in Figure 6(c-d). As x increased to 0.10 (Figure 6(e)), the dielectric constant measured at 1 and 10 kHz rapidly increased with temperature, and T_d and T_m were not visible.

Furthermore, at higher CZFMO content ($x>0.10$), the dielectric curve and loss dramatically increased with temperature, and T_d and T_m were not present at any frequency, as observed in Figure 6(g-j). These observations could be related to space charge polarization and higher ionic conductivity at higher temperatures resulting from CZFMO substitution inducing leakage currents (Nanda *et al.*, 2019).

At 1 kHz, the dielectric constant (ϵ_r) gradually decreased from 1505 to 364, while dielectric loss ($\tan\delta_r$) slightly decreased from 0.05 to between 0.02 and 0.03 as CZFMO content increased from 0 to 0.20, as listed in Table 2. A decrease in ϵ may be due to multiple factors, including poor morphology, decreased density, and increased conductivity with the addition of CZFMO. The dielectric constant (ϵ_m) and loss ($\tan\delta_m$) at T_m , measured at 1 kHz, were 6536 and 0.03, respectively, in the pure BNBT-BSN sample. However, the ϵ_m and $\tan\delta_m$ values could not be calculated as the peaks were no longer present after adding CZFMO content, implying an increase in electrical conductivity (Wu *et al.*, 2012).

Figure 7(a) and (b) illustrate the frequency-dependent ϵ and $\tan\delta$ of BNBT-BSN-xCZFM0 ceramics with varying CZFMO levels, measured from 10^2 to 10^5 Hz at room temperature. In all samples, ϵ decreased when the frequency increased, as shown in Figure 7(a). This aligns with the common understanding that the dielectric constant is highest at low frequencies because of multiple polarization mechanisms, such as electronic, space

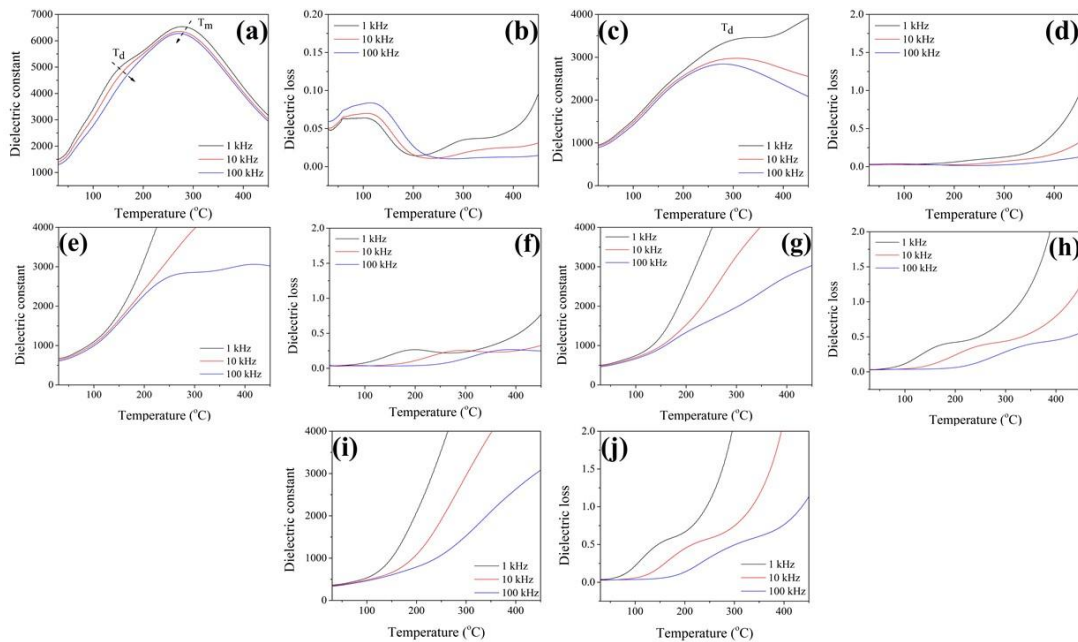


Figure 6. Temperature-dependent dielectric constant (ϵ) and loss tangent ($\tan\delta$) of BNBT-BSN-xCZFM0 ceramics at (a), (b) $x = 0$, (c), (d) $x = 0.05$, (e), (f) $x = 0.10$, (g), (h) $x = 0.15$ and (i), (j) $x = 0.20$, measured from room temperature to 500°C at frequencies of 1, 10 and 100 kHz

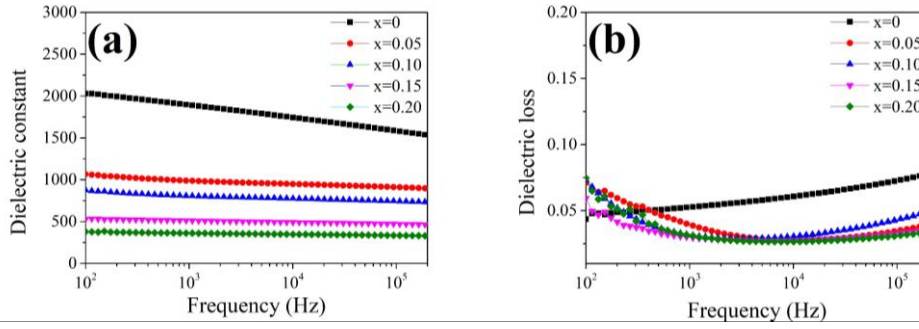


Figure 7. Frequency-dependent (a) ϵ and (b) $\tan\delta$ of BNBT-BSN-xCZFMO ceramics with various x levels, over a frequency range of 10^2 to 10^5 Hz.

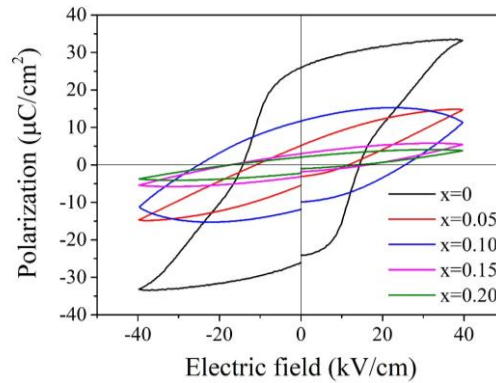


Figure 8. P-E hysteresis loops of BNBT-BSN-xCZFMO ceramics with varying x levels, under an electric field of 40 kV/cm.

charge, interfacial, ionic, and orientation polarization. As frequency increases, the dielectric constant gradually decreases until stabilizing at higher frequencies due to a reduced ability of the dipoles to align with the external electric field, causing a reduction in the orientation polarization (Zawar *et al.*, 2019). Additionally, the $\tan\delta$ at a low frequency of 10^2 Hz was similar between samples, ranging between 0.04 and 0.07, suggesting minimal impact from CZFMO content at these lower frequencies. However, as frequency increased, $\tan\delta$ tended to decrease for all CZFMO-substituted samples, as represented in Figure 7(b). This result could be explained by dipole relaxation, caused by the electric dipoles' inability to change direction at higher frequencies of the applied electric field (Singh *et al.*, 2011).

The polarization-electric field (P-E) hysteresis loops for BNBT-BSN-xCZFMO ceramics with various CZFMO content, measured under an electric field of 40 kV/cm at room temperature, are presented in Figure 8. A saturated and pinched loop was observed for the pure BNBT-BSN ceramic ($x=0$) with a high maximum polarization ($P_{\max}$ of $33.19 \mu\text{C}/\text{cm}^2$), remnant polarization (P_r of $26.09 \mu\text{C}/\text{cm}^2$), and coercive field (E_c of 14.22 kV/cm), consistent with relaxor ferroelectric behavior and

the findings of A. Luangpangai *et al.* (Luangpangai *et al.*, 2024). The P-E loop saturation decreased as CZFMO content increased. A small amount of CZFMO ($x=0.05$), led to slimmer P-E loops and a reduction in $P_{\max}$ ($14.68 \mu\text{C}/\text{cm}^2$), P_r ($5.16 \mu\text{C}/\text{cm}^2$), and E_c (11.06 kV/cm). As x increased to 0.10, the ceramic exhibited a broader P-E loop with $P_{\max}$ of $11.30 \mu\text{C}/\text{cm}^2$, P_r of $11.73 \mu\text{C}/\text{cm}^2$ and a large E_c of 24.99 kV/cm, suggesting significant leakage currents. This can be explained by decreased electrical resistance caused by the increased non-ferroelectric (ferrite) structure and the existence of secondary phases, as seen in the XRD analysis with increased CZFMO content. The added CZFMO lowered breakdown strength and voltage tolerance, resulting in a higher leakage current density (Gao *et al.* 2019).

At higher CZFMO levels ($x>0.10$), the P-E loops became narrower and $P_{\max}$, P_r and E_c decreased ($P_{\max}$ of $5.45 \mu\text{C}/\text{cm}^2$, P_r of $3.01 \mu\text{C}/\text{cm}^2$ and E_c of 17.40 kV/cm for $x=0.15$ as well as $P_{\max}$ of $3.80 \mu\text{C}/\text{cm}^2$, P_r of $2.09 \mu\text{C}/\text{cm}^2$ and E_c of 17.39 kV/cm for $x=0.20$), as listed in Table 2. These results could be attributed to the introduction of charged defect acceptor ions, such as Fe^{3+} and Co^{2+} . These ions can form electric dipoles with oxygen vacancies, functioning as pinning points, limiting

dipolar rotation, and lowering coercivity. A similar effect was found in CZFMO-doped $\text{PbZr}_{0.52}\text{Ti}_{0.48}\text{O}_3$ ceramics, as reported by A. Gupta *et al.* (Gupta *et al.*, 2011).

The magnetization-magnetic field (M-H) curves of BNBT-BSN-xCZFMO ceramics with different CZFMO content, measured over a field range of -15 to 15 kOe at room temperature, are shown in Figure 9(a-f). For the pure BNBT-BSN ceramic ($x=0$), the M-H curve had a negative slope with a remnant magnetization (M_r) of 9×10^{-5} emu/g, a saturation magnetization (M_s) of -0.0044 emu/g, and a coercive field (H_c) of 52 Oe (inset in Figure 9(a) and Table 2). The negative slope and small magnetization values suggest diamagnetic behavior. The slope became positive with a small amount of CZFMO content ($x=0.05$), indicating a transition to ferromagnetic behavior. The sample showed a higher M_r of 0.0009 emu/g, M_s of 0.0715 emu/g, and a lower H_c of 24 Oe, as seen in Figure 9(b) (inset) and Table 2. Ferromagnetic behavior develops due to structural distortion, which induces canting of the antiferromagnetic ordering chain (Futakuchi *et al.*, 2014).

As CZFMO content increased to $x=0.10$ -0.15, the M-H curves exhibited an increase in M_s (0.1411 emu/g for $x=0.10$ and 0.2485 emu/g for $x=0.15$), and a dramatic reduction in M_r (9×10^{-5} emu/g for $x=0.10$; 7×10^{-5} emu/g for $x=0.15$) and H_c (7 Oe for $x=0.10$; 3 Oe for $x=0.15$), as shown in the insets of Figure 9(c-d) and Table 2. As CZFMO increased to $x=0.20$, the M_s increased dramatically to 0.8056 emu/g with an H_c of 5 Oe (inset in Figure 9(e) and Table 2). The increase in M_s with higher CZFMO content is due to the CZFMO (ferrite) grains acting as centers of magnetization, with the overall saturation magnetization being the vector sum of these contributions (Mane *et al.*, 2018). H_c is influenced by magnetocrystalline anisotropy, microstrain, magnetic particle morphology, size distribution, shape anisotropy, and magnetic domain size (Priya *et al.*, 2015). The higher M_s and lower H_c values at higher CZFMO levels indicate the ceramic's transition to more soft magnetic behavior. However, at higher doping levels, ferromagnetic behavior deteriorated with lower M_r and H_c values, as seen in Figure 9(f) and its inset. This reduction was induced by excessive ion doping, consistent

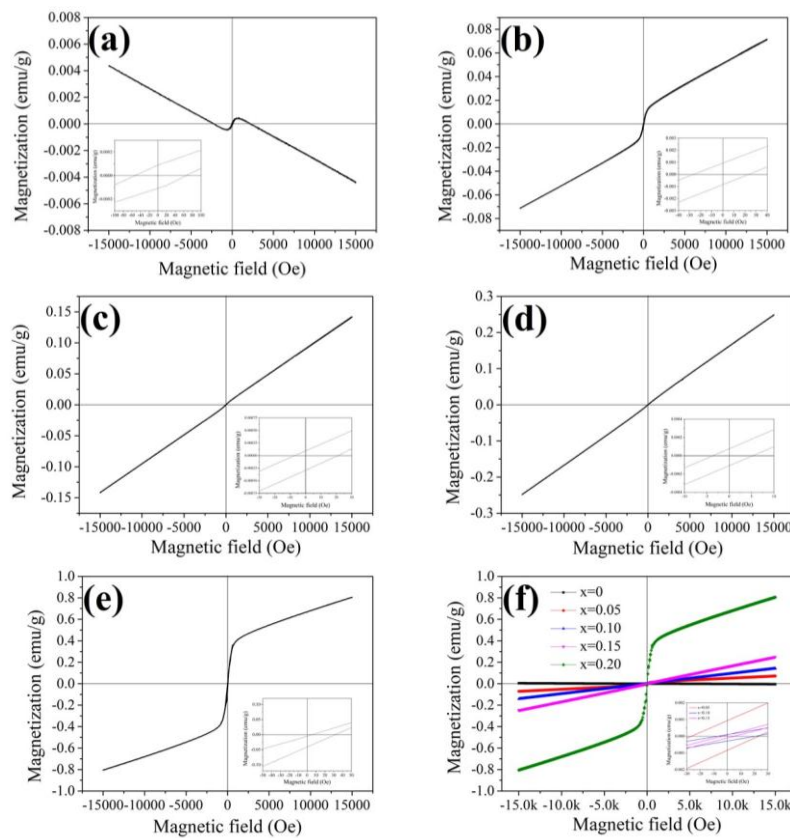


Figure 9. Magnetic behavior of BNBT-BSN-xCZFMO ceramics, indicating a transition from diamagnetic to ferromagnetic with CZFMO doping at (a) $x = 0$, (b) $x = 0.05$, (c) $x = 0.10$, (d) $x = 0.15$, (e) $x = 0.20$, and (f) all x levels

with findings reported by Y. Wan *et al.* (Wan *et al.*, 2014).

The magnetoelectric coupling coefficient (α_E) quantifies the interaction between the magnetic and electric properties of multiferroic materials and is calculated using the equation (Vopson *et al.*, 2017):

$$\alpha_E = \frac{V_{out}}{tH_{AC}} \quad (1)$$

The induced output voltage (V_{out}) was measured in a DC magnetic field with an applied AC magnetic field (H_{AC}) of 10 Oe and a frequency of 1.5 kHz, while t is the sample thickness. The α_E values for BNBT-BSN- x CZFMO ceramics with various CZFMO levels with a DC magnetic field (H_{dc}) ranging from 0 to 5000 Oe are shown in Figure 10. At H_{dc} of 1000 Oe, α_E values continually increased from 1.317 to 5.111 mV/Oe·cm when x increased from 0 to 0.20, as listed in Table 2. The result could be related to the higher magnetostriction produced by the higher mass ratio of CZFMO contents, consistent with a similar report by C. S. Chitra Lekha *et al.* (Chitra Lekha *et al.*, 2019).

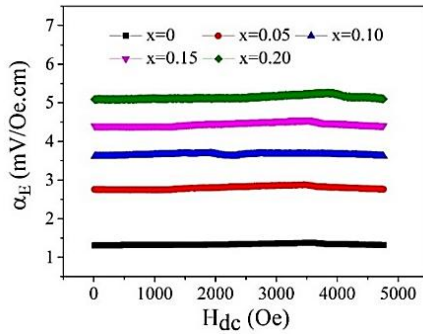


Figure 10. Variations in α_E and H_{dc} of BNBT-BSN- x CZFMO ceramics with different x levels

The 0.90BNBT-BSN-0.10CZFMO ceramic, fabricated using the solid-state combustion technique, exhibited optimal electrical, magnetic and magnetoelectric properties (ϵ_r of 674, P_{max} of 11.30 μ C/cm², P_r of 11.73 μ C/cm², M_s of 0.1411 emu/g, M_r of 9×10^{-4} emu/g, H_c of 7 Oe and α_E of 3.655 mV/Oe·cm). Other multiferroic ceramics, such as 0.90Bi_{0.465}Na_{0.465}Ba_{0.07}TiO₃-0.10Ni_{0.3}Cu_{0.08}Zn_{0.62}Fe₂O₄, exhibited a P_{max} of ~ 1 μ C/cm², P_r of ~ 0.7 μ C/cm² and M_r of 7.55×10^{-6} emu/g (Nanda *et al.*, 2019). Similarly, 0.90Ba_{0.9}Sr_{0.1}Zr_{0.04}Ti_{0.96}O₃-0.10Ni_{0.8}Zn_{0.2}Fe₂O₄ multiferroic ceramics had a P_r of ~ 8 μ C/cm² and α_E of 1.8 mV/Oe·cm (Rani *et al.*, 2017), while 0.7BaTiO₃-0.3Ni_{0.7}Zn_{0.3}Fe₂O₄ ceramics had an α_E of 0.124 mV/Oe·cm (Zhang *et al.*, 2013). These comparisons highlight that 0.90BNBT-BSN-

0.10CZFMO ceramics are a promising candidate for application in magnetoelectric devices and materials.

Conclusions

In summary, (1- x)[0.99Bi_{0.47}Na_{0.47}Ba_{0.06}TiO₃-0.01Ba(Sn_{0.70}Nb_{0.24})O₃]- x Co_{0.6}Zn_{0.4}Fe_{1.7}Mn_{0.3}O₄ multiferroic ceramics (abbreviated as BNBT-BSN- x CZFMO) with $x=0$ -0.20 were successfully synthesized using the solid-state combustion technique. The phase structure, microstructure, electrical, magnetic, and magnetoelectric properties were studied. The coexistence of the perovskite and spinel phases was confirmed with CZFMO substitution, with a reduction in lattice parameters as CZFMO content increased. In addition, ferroelectric and ferrite grains were detected in the microstructure of BNBT-BSN ceramics when CZFMO was added. The porosity increased while the measured density decreased with higher CZFMO contents. The ϵ_r dropped with increasing CZFMO, caused by inferior microstructure and density, while ϵ_m was unmeasurable due to high electrical conductivity. Due to a high leakage current, the P-E loops became slimmer at $x=0.05$ but broadened significantly at $x=0.10$, with higher P_{max} , P_r , and E_c . Further additions of CZFMO content caused the P-E loops to become thinner and reduced the P_{max} , P_r , and E_c . The M-H curves showed a transition to ferromagnetic behavior with increasing CZFMO, as M_s increased while H_c decreased, implying enhanced soft magnetic behavior. The α_E values gradually increased from 1.317 to 5.111 mV/Oe·cm when x increased from 0 to 0.20. The study demonstrates improved magnetoelectric properties at optimal doping levels, suggesting potential applications in multiferroic devices.

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